Nonlocal Wasserstein distance and the associated gradient flows

Dejan Slepčev Carnegie Mellon University

Simons Institute

October 25, 2021.

Collaborations and References

- Nonlocal Wasserstein distance with Andrew Warren (in preparation)
- Nonlocal-interaction equation on graphs: gradient flow structure and continuum limit with Antonio Esposito, Francesco Patacchini, and André Schlichting, ARMA 2021
- Clustering dynamics on graphs: from spectral clustering to mean shift through Fokker-Planck interpolation with Katy Craig and Nicolás García Trillos
- *Gradient flows of the entropy for jump processes*, Erbar 2014, Annales IHP Prob and Stat.
- Gromov–Hausdorff limit of Wasserstein spaces on point clouds García Trillos 2020 Calc. Var. and PDE

Consider probability measures with finite second moment $\rho_0, \rho_1 \in \mathcal{P}_2(\mathbb{R}^d)$ Let \mathcal{A} be the set of paths from ρ_0 to ρ_1 :

$$\mathcal{A} = \{ (\rho, \mathbf{v}) : \rho(\cdot, t) \in L^1(\mathbb{R}^d, [0, \infty)), \mathbf{v}_t \in L^2(d\mu_t), \\ \partial_t \rho + \operatorname{div}(\rho \, \mathbf{v}) = \mathbf{0} \text{ on } \mathbb{R}^d \times [0, t] \\ \rho(\cdot, 0) = \rho_0 \text{ and } \rho(\cdot, 1) = \rho_1 \}$$

Theorem

$$d_W^2(\rho_0,\rho_1) = \inf_{(\rho,\nu)\in\mathcal{A}} \int_0^1 \int |\boldsymbol{\nu}(\boldsymbol{x},\boldsymbol{t})|^2 d\rho_t(\boldsymbol{x}) dt$$

Nonlocal calculus

Let $\eta : \mathbb{R}^d \times \mathbb{R}^d \to [0, \infty]$ describe the nonlocal transportation connections. Assume $\eta(x, y) = \eta(y, x)$. Let $G = \{(x, y) : \eta(x, y) > 0\}$.

Nonlocal Gradient. Given $\varphi : \mathbb{R}^d \to \mathbb{R}$, the gradient $\overline{\nabla} \varphi : G \to \mathbb{R}$

$$\overline{\nabla}\varphi(\mathbf{x},\mathbf{y})=\varphi(\mathbf{y})-\varphi(\mathbf{x}).$$

Nonlocal Divergence. Given $j \in \mathcal{M}(G)$ – signed measures on G

$$\overline{\nabla} \cdot \mathbf{j}(\mathbf{x}) = \int \eta(\mathbf{x}, \mathbf{y}) \mathbf{j}(\mathbf{x}, d\mathbf{y}).$$

For any $\mathbf{j} \in \mathcal{M}(G)$, its *nonlocal divergence* $\overline{\nabla} \cdot \mathbf{j} \in \mathcal{M}(\mathbb{R}^d)$ is defined as η -weighted adjoint of $\overline{\nabla}$, i.e.,

$$\int \phi \, d\overline{\nabla} \cdot \mathbf{j} = -\frac{1}{2} \iint_{G} \overline{\nabla} \phi(\mathbf{x}, \mathbf{y}) \eta(\mathbf{x}, \mathbf{y}) \, d\mathbf{j}(\mathbf{x}, \mathbf{y}).$$

Nonlocal continuity equation

What is the nonlocal analog of the continuity equation:

$$\partial_t \rho_t + \nabla \cdot \mathbf{j}_t = 0$$
 with flux $\mathbf{j}_t(\mathbf{x}) = \rho_t(\mathbf{x})\mathbf{v}_t(\mathbf{x})$?

Nonlocal fluxes j_t are defined on edges $(x, y) \in G$ while the densities are defined at points

$$\partial_t \rho_t(\mathbf{x}) + (\overline{\nabla} \cdot \mathbf{j}_t)(\mathbf{x}) = \partial_t \rho_t(\mathbf{x}) + \int_G \eta(\mathbf{x}, \mathbf{y}) \mathbf{j}_t(\mathbf{x}, d\mathbf{y}) = \mathbf{0} \ .$$

What is the nonlocal relationship between **j** and velocity v, given ρ_t ?

Problem: There is no canonical way to define density along edges. *[Maas '11], [Mielke '11], [Chow, Huang, Li, Zhou '12], [Erbar '14]* use averaging functions $\theta : \mathbb{R}_+ \times \mathbb{R}_+ \to \mathbb{R}_+$:

$$\mathbf{j}(\mathbf{x},\mathbf{y}) = \theta(\rho(\mathbf{x}),\rho(\mathbf{y})) \mathbf{v}(\mathbf{x},\mathbf{y}).$$

For any $r, s \ge 0$, $\theta(r, r) = r$, θ is increasing in r and s.

Nonlocal Wasserstein distance

Erbar '14, generalizes Maas '11

Given $\rho_0, \rho_1 \in \mathcal{P}_2(\mathbb{R}^d)$, consider vector fields $v : [0, 1] \times G \to \mathbb{R}$ such that the solution of the continuity equation

$$\partial_t \rho(\mathbf{x}) + \overline{\nabla} \cdot (\theta(\rho(\mathbf{x}), \rho(\mathbf{y})) \mathbf{v}(\mathbf{x}, \mathbf{y})) = 0$$

with $\rho(0) = \rho_0$ and $\rho(1) = \rho_1$.

Admissible paths, $CE(\rho_0, \rho_1)$, are all of the solutions $(\rho_t, v_t)_{t \in [0,1]}$ generated by above vector fields.

$$\mathcal{W}^2_{\eta}(\rho_0,\rho_1) := \frac{1}{2} \inf_{\mathsf{CE}(\rho_0,\rho_1)} \int_0^1 \int_G |v_t(x,y)|^2 |\theta(\rho_t(x),\rho_t(y))\eta(x,y) dx dy dt$$

On graphs

$$\mathcal{W}_{\eta}^{2}(\rho_{0},\rho_{1}) := \frac{1}{2} \inf_{\mathsf{CE}(\rho_{0},\rho_{1})} \int_{0}^{1} \sum_{x \in V} \sum_{y \in V} |v_{t}(x,y)|^{2} |\theta(\rho_{t}(x),\rho_{t}(y))\eta(x,y)dt$$

- $\theta(r, 0) = 0$, which includes geometric and logarithmic mean [Erbar, Maas]
- θ(r, 0) > 0, which includes arithmetic mean. The set of tangent fluxes needs to be restricted to a cone.

• $\theta(r, s, j) = \begin{cases} r & \text{if } j > 0 \\ s & \text{if } j \le 0 \end{cases}$ is the upwind interpolation. The resulting "distance" is not symmetric.

Graph Laplacian

•
$$V_n = \{x_1, \ldots, x_n\}$$
, similarity matrix *W*:

$$W_{ij} := \eta \left(|\mathbf{x}_i - \mathbf{x}_j| \right).$$

The weighted degree of a vertex is $d_i = \sum_j W_{i,j}$.

Graph Laplacian

$$L = D - W$$
,

where
$$D = \text{diag}(d_1, \ldots, d_n)$$
.

• Graph heat equation

$$\frac{d}{dt}\rho(\mathbf{x}_i) = -Lu(\mathbf{x}_i).$$

Graph heat equation as graph Wasserstein gradient flows

[Maas '11], [Mielke '11], [Chow, Huang, Li, Zhou '12] The graph heat equation

$$\frac{d}{dt}
ho = -L
ho$$

is the gradient of entropy

$$\Xi(\rho) = \sum_{x \in V} \rho(x) \ln \rho(x)$$

with respect to the graph Wasserstein distance $d_{w,G}^2$ corresponding to

$$\theta(r,s) = rac{r-s}{\ln r - \ln s}$$
 and $\theta(r,0) = 0$

For $F(\rho) = \sum_{x \in V} U(x)\rho(x)$ where *U* is a smooth function, and $\rho(\cdot, 0) = \delta_{x_1}$ the gradient flow is $\partial_t \rho(x) = 0$ for all *x*. The support of the solution cannot expand.

Graph heat equation as graph Wasserstein gradient flows

[Maas '11], [Mielke '11], [Chow, Huang, Li, Zhou '12] The graph heat equation

$$\frac{d}{dt}
ho = -L
ho$$

is the gradient of entropy

$$\Xi(\rho) = \sum_{x \in V} \rho(x) \ln \rho(x)$$

with respect to the graph Wasserstein distance $d_{w,G}^2$ corresponding to

$$\theta(r,s) = rac{r-s}{\ln r - \ln s}$$
 and $\theta(r,0) = 0$

For $F(\rho) = \sum_{x \in V} U(x)\rho(x)$ where *U* is a smooth function, and $\rho(\cdot, 0) = \delta_{x_1}$ the gradient flow is $\partial_t \rho(x) = 0$ for all *x*. The support of the solution cannot expand.

Fractional Heat Equation

Nonlocal diffusion equation

$$\partial_t \rho = \int (\rho(\mathbf{y}) - \rho(\mathbf{x})) \eta(\mathbf{x} - \mathbf{y}) d\mathbf{y}$$

is the gradient flow of entropy

$$E(
ho) = \int
ho(x) \ln
ho(x) dx$$

with respect to the gnonlocal Wasserstein distance \mathcal{W}_η corresponding to

$$\theta(r,s) = rac{r-s}{\ln r - \ln s} \text{ and } \theta(r,0) = 0$$

In particular, as Erbar observed, for $s \in (0, 2)$ the fractional heat equation

$$\partial_t \rho = \int \frac{1}{|x-y|^{-d-s}} (\rho(y) - \rho(x)) \, dy = \Delta^{s/2} \rho$$

is the gradient flow of entropy with respect to the nonlocal Wasserstein distance for $\eta(z) = \frac{1}{|z|^{-d-s}}$.

Difficulties:

- ρ may contain atoms
 - \Rightarrow measure valued framework
- Benamou-Brenier functional is not jointly convex in (ρ_t, ν_t) ⇒ flux variables
- Upwind metric is is only positively homogeneous: $g(v, v) \neq g(-v, -v)$. The geometric structure is Finslerian rather than Riemannian.
- In the general framework the underlying space ρ is supported within is described by measure μ . For most of the talk μ is the Lebesgue measure.

Nonlocal continuity equation in measure valued flux form

A pair $(\rho_t, \mathbf{j}_t)_{t \in [0, T]} \in CE_T$ provided that $(\rho_t, \mathbf{j}_t) \in \mathcal{P}(\Omega) \times \mathcal{M}(G)$ for all $t \in [0, T]$:

$$\partial_t \rho_t + \overline{\nabla} \cdot \mathbf{j}_t = 0$$
 in $C^{\infty}_{c}([0, T) \times \Omega)^*$

That is for smooth test functions φ

$$\int_0^T \int_\Omega \partial_t \varphi_t(x) d\rho_t(x) dt + \int_0^T \iint_G \overline{\nabla} \varphi_t(x, y) \eta(x, y) d\mathbf{j}_t(x, y) dt = 0.$$

Nonlocal continuity equation and action

Nonlocal continuity equation in measure valued flux form

A pair $(\rho_t, \mathbf{j}_t)_{t \in [0, T]} \in CE_T$ provided that $(\rho_t, \mathbf{j}_t) \in \mathcal{P}(\Omega) \times \mathcal{M}(G)$ for all $t \in [0, T]$:

$$\partial_t \rho_t + \nabla \cdot \mathbf{j}_t = 0$$
 in $C^{\infty}_c([0, T) \times \Omega)^*$

Action [for upwind flux]

For $\mathbf{j} \in \mathcal{M}(G)$, set $|\lambda| = |\rho \otimes \mu| + |\mu \otimes \rho| + |\mathbf{j}| \in \mathcal{M}^+(G)$ and define

$$\mathcal{A}(\mu;\rho,\boldsymbol{j}) = \iint_{\boldsymbol{G}} \left(\alpha \left(\frac{d\boldsymbol{j}}{d|\lambda|}, \frac{d(\rho \otimes \mu)}{d|\lambda|} \right) + \alpha \left(-\frac{d\boldsymbol{j}}{d|\lambda|}, \frac{d(\mu \otimes \rho)}{d|\lambda|} \right) \right) \, \eta d|\lambda|.$$

where, the convex, and pos. one-homogeneous function $\boldsymbol{\alpha}$ is defined by

$$\alpha(j,r) := \begin{cases} \frac{(j_+)^2}{r} & \text{if } r > 0, \qquad \text{with } j_+ = \max\{0, j\} \\ 0 & \text{if } j \le 0 \text{ and } r = 0, \\ \infty & \text{if } j > 0 \text{ and } r = 0. \end{cases}$$

Finite action leads to upwind flux

Proposition

Let $(\rho, \mathbf{j}) \in \mathcal{P}(\Omega) \times \mathcal{M}(\Omega)$ such that $\mathcal{A}(\mu; \rho, \mathbf{j}) < \infty$, then:

• there exists a measurable nonlocal vector field $v: G \to \mathbb{R}$ such that

$$d\,\mathbf{j}(x,y) = \mathbf{v}(x,y)_{+} \, d\,\rho(x) \, d\,\mu(y) - \mathbf{v}(x,y)_{-} \, d\,\mu(x) \, d\,\rho(y) \quad \text{and} \\ \mathcal{A}(\mu;\rho,\mathbf{j}) = \iint_{\mathcal{G}} \left(|\mathbf{v}(x,y)_{+}|^{2} + |\mathbf{v}(y,x)_{-}|^{2} \right) \eta(x,y) d\rho(x) d\mu(y) \, .$$

• there exists an antisymmetric $\textit{\textbf{j}}^{as} \in \mathcal{M}^{as}(G)$ such that

$$\overline{\nabla} \cdot \boldsymbol{j} = \overline{\nabla} \cdot \boldsymbol{j}^{\boldsymbol{as}}, \quad ext{that is} \quad \iint_{\boldsymbol{G}} \overline{\nabla} \phi \, \eta \, \boldsymbol{d} \, \boldsymbol{j} = \iint_{\boldsymbol{G}} \overline{\nabla} \phi \, \eta \, \boldsymbol{d} \, \boldsymbol{j}^{\boldsymbol{as}} \quad \forall \phi \in \boldsymbol{C}^{\infty}_{\boldsymbol{c}}(\Omega)$$

and an antisymmetric $v^{as}: G \to \mathbb{R}$ with

$$\mathcal{A}(\mu;\rho,\boldsymbol{j}^{as}) = 2 \iint_{G} |\boldsymbol{v}^{as}(\boldsymbol{x},\boldsymbol{y})_{+}|^{2} \eta d\rho(\boldsymbol{x}) d\mu(\boldsymbol{y}) \leq \mathcal{A}(\mu;\rho,\boldsymbol{j}).$$

Lower semicontinuity and compactness

Lower semicontinuity

if $\mu^n \rightarrow \mu$ in $\mathcal{M}(\Omega)$, $\rho^n \rightarrow \rho$ in $\mathcal{P}(\Omega)$, and $\mathbf{j}^n \rightarrow \mathbf{j}$ in $\mathcal{M}(\mathbf{G})$, then

$$\liminf_{n\to\infty} \mathcal{A}(\mu^n;\rho^n,\boldsymbol{j}^n) \geq \mathcal{A}(\mu;\rho,\boldsymbol{j})$$

Compactness

Let $(\rho^n, \mathbf{j}^n) \in CE_T$ for each $n \in \mathbb{N}$ such that $\sup_n \int_0^T \mathcal{A}(\rho_t^n, \mathbf{j}_t^n) dt < \infty$. Then, there exists $(\rho, \mathbf{j}) \in CE_T$ such that

$$\rho_t^n \rightarrow \rho_t \quad \text{in } \mathcal{P}_2(\Omega) \text{ for all } t \in [0, T]$$
$$\mathbf{j}^n \rightarrow \mathbf{j} \quad \text{in } \mathcal{M}_{\text{loc}}(\mathbf{G} \times [0, T]).$$

Moreover $\liminf_{n\to\infty} \int_0^T \mathcal{A}(\rho_t^n, \mathbf{j}_t^n) \, dt \geq \int_0^T \mathcal{A}(\rho_t, \mathbf{j}_t) \, dt.$

Compactness of solutions to CE

Assumption (weight function)

The μ -measurable nonnegative symmetric lsc. function $\eta: G \to \mathbb{R}$ satisfies:

• The measure $\eta(\cdot,\cdot)d\mu$ is uniformly integrable close to diagonal, that is

$$\lim_{\varepsilon\to 0}\sup_{x\in\Omega}\int_{B_{\varepsilon}(x)}|x-y|^2\,\eta(x,y)d\mu(y)=0,\quad B_{\varepsilon}(x)=\big\{y\in\Omega:|x-y|<\varepsilon\big\}.$$

Compactness: Let $(\rho^n, j^n) \in CE_T$ for each $n \in \mathbb{N}$ such that $\sup_{n \in \mathbb{N}} M_2(\rho_0^n) < \infty$ and $\sup_n \int_0^T \mathcal{A}(\rho_t^n, j_t^n) dt < \infty$. Then, there exists $(\rho, j) \in CE_T$ such that

$$\begin{array}{ll} \rho_t^n \rightharpoonup \rho_t & \text{in } \mathcal{P}_2(\Omega) \text{ for all } t \in [0,T] \\ \boldsymbol{j}^n \rightharpoonup \boldsymbol{j} & \text{in } \mathcal{M}_{\mathrm{loc}}(\boldsymbol{G} \times [0,T]). \end{array}$$

Moreover, the action is lower semicontinuous

$$\liminf_{n\to\infty}\int_0^T \mathcal{A}(\rho_t^n, \boldsymbol{j}_t^n)\,d\,t\geq\int_0^T \mathcal{A}(\rho_t, \boldsymbol{j}_t)\,d\,t.$$

Definition

For $\rho_0, \rho_1 \in \mathcal{P}_2(\Omega)$ the nonlocal upwind Wasserstein quasimetric

$$\mathcal{W}_{\eta}(\rho_0,\rho_1)^2 = \inf\left\{\int_0^1 \mathcal{A}(\rho_t,\boldsymbol{j}_t)dt : (\rho,\boldsymbol{j}) \in \mathsf{CE}(\rho_0,\rho_1)\right\}$$

Properties:

- Minimum is attained for $(\rho, \mathbf{j}) \in CE(\rho_0, \rho_1)$ with $\mathcal{A}(\rho_t, \mathbf{j}_t) = W_{\eta}(\rho_0, \rho_1)^2$.
- \mathcal{W}_{η} is jointly narrowly lower semicontinuous.
- For upwind interpolation we will use T_{η} instead of W_{η} . Note that $T\eta$ is not symmetric.

Expell cost

Proposition [Warren and S.]

• If η is nonintegrable: $\eta(x, y) = \eta(|x - y|)$ and $\eta(r) > r^{-d-s}$ when $r \leq \delta$, then

$$\mathcal{W}_\eta(\delta_{\mathbf{0}}, \boldsymbol{c}\chi_{\boldsymbol{B}(\mathbf{0},\delta)}) \lesssim \delta^{\boldsymbol{s}/2}.$$

If η(x, y) = η(x − y) and ∫_{ℝ^d} ηdx < ∞ and we consider arithmetic mean or upwind interpolation then there exists c > 0 such that for all ν⊥ρ

$$\mathcal{W}_{\eta}(\delta_{0},\nu)\geq c.$$

• If $\eta(x, y) = \eta(x - y)$ and $\int_{\mathbb{R}^d} \eta dx < \infty$ and we consider interpolation with $\theta(r, 0) = 0$ then $\nu \perp \rho$

$$\mathcal{W}_{\eta}(\delta_0,\nu) = \infty.$$

Furthermore $\mathcal{T}_{\eta}(\nu, \delta_0) = \infty$.

Topology

Erbar showed that the topology metrized by NLW is at least as strong as the the one generated by Wasserstein metric (narrow convergence plus moment control)

Proposition [Warren and S.]

- If η is nonintegrable: $\eta(x, y) = \eta(|x y|)$ and $\eta(r) > r^{-d-s}$ when $r \leq \delta$, then the topology generated by W_{η} on $\mathcal{P}((B(0, R)))$ is the weak topology.
- If η is integrable: If η(x, y) = η(x − y) and ∫_{ℝ^d} ηdx < ∞ and we consider arithmetic mean then there exist 0 < c₁, c₂ such that for all ρ₀, ρ₁ ∈ P((B(0, R))

$$\boldsymbol{c}_1 \| \rho_1 - \rho_0 \|_{\mathcal{T}V} \leq \mathcal{W}(\rho_0, \rho_1) \leq \boldsymbol{c}_2 \left(\| \rho_1 - \rho_0 \|_{\mathcal{T}V} + \boldsymbol{R} \sqrt{\boldsymbol{d}_{\textit{Monge}}(\rho_0, \rho_1)} \right)$$

Theorem [Warren and S.]

Upper bound. In the cases where the expel cost is finite (and not including upwind interpolation) and η compactly supported

$$arepsilon \mathcal{W}_{arepsilon,\eta}(\mu_0,\mu_1) \leq \left(rac{1}{\sqrt{\sigma_\eta}}
ight) d_W(\mu_0,\mu_1) + \mathcal{O}(\sqrt{arepsilon}).$$

where $\eta_{\varepsilon}(z) = \frac{1}{\varepsilon^d} \eta \frac{z}{\varepsilon}$, ω_d is the volume of the unit ball, and

$$\sigma_{\eta} = \frac{1}{d} \int_{\mathbb{R}^d} |x|^2 \eta(x) dx.$$

Lower bound. Assume both $\mu_0, \mu_1 \in \mathcal{P}(\mathbb{R}^d)$ are supported inside $\overline{B}(0, R)$ with $R \geq 1$ and η compactly supported. Then,

$$d_W^2(\mu_0,\mu_1) \leq \varepsilon^2 \sigma_\eta \mathcal{W}_{\eta,\varepsilon}^2(\mu_0,\mu_1) + C R^2 \sqrt{\varepsilon}.$$

Upper bound: elements of the proof I

We use the Wasserstein geodesic to build a competitor. This includes two levels of smoothing.

1. Exact solutions to nonlocal transport. Let

$$\zeta(r) = \int_r^\infty s\eta(s)ds.$$

Consider a solution of the continuity equation

$$\partial_t \rho + \nabla \cdot J = 0.$$
Let $\rho_{\zeta} = \rho * \zeta$ and $J_{\zeta} = J * \zeta$. Then $\partial_t \rho_{\zeta} + \nabla \cdot (J_{\zeta}) = 0$.
Let $j(x, y) = (y - x) \cdot J(y)$. Then

$$\partial_t \rho_{\zeta} + \overline{\nabla} \cdot j = \partial_t \rho_{\zeta} + \int j(x, y) \eta(x - y) dy = 0.$$

We use the Wasserstein geodesic to build a competitor. This includes two levels of smoothing.

2. Smoothing that controls the interpolation. The problem arises if $\rho(x) - \theta(\rho(x), \rho(y)) > c > 0$. Let $K(x) = c e^{-|x|}$ where $\frac{1}{c} = \int_{\mathbb{R}^d} e^{-|x|} dx$. Let $K_{\delta}(x) = \frac{1}{\delta^d} K\left(\frac{x}{\delta}\right)$. Consider $\mu \in \mathcal{P}_2(\mathbb{R}^d)$. Let $\mu_{\delta} = \mu * K_{\delta}$ If $|y - x| < \delta$, $\mu_{\delta}(y) \le \mu_{\delta}(x) \left(1 + \frac{3}{\delta}|y - x|\right)$. We use the dual formulation to provide competitor. In particular the nonlocal Hamilton-Jacobi equation. For graphs developed in *Gangbo, Li, Mou '19, Erbar, Maas, Wirth '20*.

Background:

Lemma. Suppose that μ_0 and μ_1 are probability measures supported within B(0, R). Then,

$$\frac{1}{2}d_W^2(\mu_0,\mu_1) = \sup_{\phi_t \in \mathcal{BL}([0,1] \times \mathbb{R}^d)} \left\{ \int \phi_1 d\mu_1 - \int \phi_0 d\mu_0 : \partial_t \phi_t + \frac{1}{2} |\nabla \phi_t|^2 \leq 0 \right\}$$

it holds that the optimal Hamilton-Jacobi subsolution has the property that $Lip(\phi_t) \leq 2R$, for Lebesgue-almost all $t \in [0, 1]$.

A Lipschitz function $\phi : \mathbb{R}^d \times [0, 1] \to \mathbb{R}$ is a nonlocal Hamilton-Jacobi subsolution, $\phi_t \in HJ_{NL}^1$ if, for a.e. $t \in (0, T)$ for all probability measures $\mu \in \mathcal{P}_2(\mathbb{R}^d)$, and for any λ such that $Leb \ll \lambda$,

$$\int \partial_t \phi_t d\mu + \frac{1}{2} \int (\phi_t(y) - \phi_t(x))^2 \theta\left(\frac{d\mu}{d\lambda}(x), \frac{d\mu}{d\lambda}(y)\right) \eta_\varepsilon(x, y) d\lambda(x) d\lambda(y) \leq 0.$$

Then, the duality formula we expect to hold is

$$\frac{1}{2}\mathcal{W}_{\eta,\varepsilon}^2(\mu_0,\mu_1) = \sup\left\{\int \phi_1(x)d\mu_1(x) - \int \phi_0(x)\mu_0(x) : \phi_t \in \mathsf{HJ}_{\mathsf{NL}}^1\right\}.$$

For technical reasons, we introduce a "smoothed version" of the nonlocal Wasserstein distance instead.